

Open heavy flavour nuclear modification factor in ALICE

André Mischke* for the ALICE Collaboration

ERC-Research Group QGP-ALICE, Utrecht University

Princetonplein 5, 3584 CC Utrecht, the Netherlands

E-mail: a.mischke@uu.nl

The ALICE experiment has measured charm and beauty production in pp and Pb–Pb collisions at $\sqrt{s} = 2.76$ and 7 TeV and $\sqrt{s_{NN}} = 2.76$ TeV, respectively, via the exclusive reconstruction of hadronic D meson decays and semi-leptonic D and B meson decays. In this contribution, recent results on the nuclear modification factor for open charmed mesons and heavy-flavour decay electrons at central rapidity and muons at forward rapidity are presented and discussed.

50th International Winter Meeting on Nuclear Physics

23-27 January 2012

Bormio, Italy

*Speaker.

1. Introduction

High-energy collisions of heavy atomic nuclei allow to explore the behaviour of strongly interacting matter at high temperatures, where a new phase of matter, the Quark-Gluon Plasma (QGP), is predicted to exist (for a recent review see [1]). Heavy quarks (charm and beauty), abundantly produced in heavy ion collisions at the Large Hadron Collider (LHC), provide sensitive penetrating probes of such matter. Due to their large mass, they are believed to be predominantly produced in the initial state of the collision by gluon fusion processes and thus provide information about the hottest initial phase. These heavy quarks propagate through the hot and dense QCD matter and lose energy through medium-induced gluon radiation (colour charge dependent) and collisions with the medium. Moreover, theoretical models predicted that heavy quarks should experience smaller energy loss than light quarks due to the suppression of small angle gluon radiation (*dead-cone effect* [2, 3]). Thus, the study of heavy-flavour production in nucleus-nucleus collisions provides key tests of parton energy-loss models yielding profound insight into the properties of the produced QCD matter.

Medium effects are typically quantified using the nuclear modification factor R_{AA} where the particle yield in Pb–Pb collisions is divided by the yield in pp reactions scaled by the number of binary collisions. $R_{AA}=1$ would indicate that no nuclear effects, such as Cronin effect, shadowing or gluon saturation, are present and that nucleus-nucleus collisions can be considered as an incoherent superposition of nucleon-nucleon interactions. By comparing the nuclear modification factor of charged pions ($R_{AA}^{\pi^{\pm}}$), mostly originating from gluon fragmentation at this collision energy, with that of hadrons with charm R_{AA}^D and beauty R_{AA}^B the dependence of the energy loss on the parton nature (quark/gluon) and mass can be investigated. A mass ordering pattern $R_{AA}^{\pi^{\pm}} < R_{AA}^D < R_{AA}^B$ is expected [4].

2. ALICE detector setup, trigger and data set

ALICE, A Large Ion Collider Experiment, is the dedicated detector for measurements on high energy heavy-ion collisions [5, 6]. Results with the first LHC Pb–Pb data are reported in [7, 8]. Its characteristic features are the very low momentum cut-off (100 MeV/c), the low material budget and the excellent particle identification and vertexing capabilities in a high multiplicity environment. At mid-rapidity, tracking with up to 159 three-dimensional space-points and particle identification through the measurement of the specific ionisation energy loss (dE/dx) is performed using the large volume Time Projection Chamber (TPC), located inside the large solenoidal magnet with a field of $B = 0.5$ T. The TPC has a coverage from -0.9 to 0.9 in pseudo-rapidity and 2π in azimuth. The relative dE/dx resolution is 5-6% and the relative momentum resolution $\Delta p/p$ is 5% at 100 GeV/c. The particle identification up to about $p_t < 2$ GeV/c is performed with the Time Of Flight (TOF) detector, which has an intrinsic time resolution better than 100 ps. The Inner Tracking System consists of six concentric cylindrical layers of silicon detectors and provides excellent reconstruction of displaced vertices with a transverse impact parameter resolution better than 65 μm for $p_t > 1$ GeV/c [9]. The track impact parameter resolution in the transverse plane in pp and Pb–Pb collisions is depicted in Fig. 1. The total material budget for radial tracks in the transverse plane is $\approx 7.7\%$ of the radiation length. At forward rapidity ($-4 < \eta < -2.5$) muons are identified

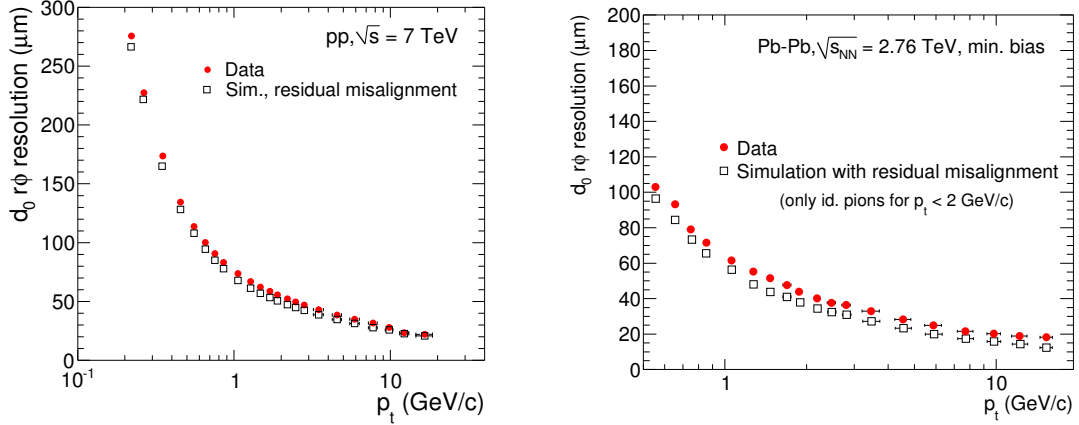


Figure 1: Track impact parameter (d_0) resolution in the transverse plane ($r\phi$ direction) as a function of p_t in pp (left panel) and Pb-Pb collisions (right panel) comparing data and Monte Carlo simulation [9]. This resolution includes the uncertainty in the primary vertex position, which is reconstructed excluding the track being probed. For $p_t < 2$ GeV/c, pion identification by the TPC or TOF detectors is required.

in the muon spectrometer. The muon spectrometer detects muons with momentum larger than 4 GeV/c and is composed of two absorbers, a dipole magnet providing a field integral of 3 Tm, and tracking and trigger chambers. Tracking is performed by means of five tracking chambers with an intrinsic spatial resolution better than 100 μm .

The Pb-Pb data presented in this paper were collected with a minimum-bias trigger in November-December 2010 during the first run with heavy-ions at the LHC at a centre-of-mass energy of $\sqrt{s_{NN}} = 2.76$ TeV per nucleon-nucleon pair. This trigger is based on the information of the two innermost silicon layers ($|\eta| < 2$) and the VZERO scintillator hodoscopes ($2.8 < \eta < 5.1$ and $-3.7 < \eta < -1.7$). The efficiency for triggering hadronic interactions was 100% for Pb-Pb collisions in the centrality range considered in the analyses. Beam background collisions were removed offline on the basis of the timing information provided by the VZERO and the Zero Degree Calorimeters (located near the beam pipe at $z \pm 114$ m from the interaction point). Only events with a vertex position within 10 cm from the centre of the detector along the beam line were used, resulting in 17M events. The forward single-muon arm triggers in coincidence with the minimum-bias trigger. The Pb-Pb data are classified based on their centrality defined in terms of percentiles of the hadronic Pb-Pb cross section and determined from the distribution of the summed amplitudes in the VZERO detector. This distribution was fitted using the Glauber model for the geometrical description of the nuclear collision together with a two-component model for particle production [10]. The pp results correspond to 315M and 180M minimum-bias events at $\sqrt{s} = 7$ TeV for the D meson (integrated luminosity: 5 nb^{-1}) and single electron analyses (2.6 nb^{-1}), respectively, and 2M muon triggered events (16.5 nb^{-1}).

3. Results

The following sections give a brief description of the analyses and results on charm and beauty production in pp and Pb-Pb collisions at $\sqrt{s} = 2.76$ and 7 TeV and $\sqrt{s_{NN}} = 2.76$ TeV, respectively,

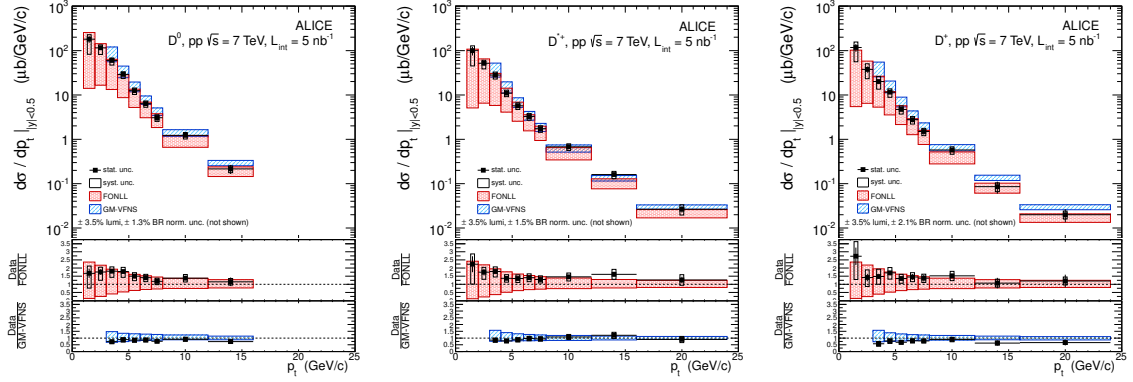


Figure 2: (colour online) p_t -differential inclusive cross section for prompt D^0 , D^{*+} and D^+ mesons for $|\eta| < 0.5$ in pp collisions at $\sqrt{s} = 7$ TeV [11], compared with FONLL [14] and GM-VFNS [15] theoretical calculations. The symbols are positioned horizontally at the centre of each p_t interval. The normalization uncertainty is not shown (3.5% from the minimum-bias cross section plus the branching ratio uncertainties).

via the exclusive reconstruction of hadronic and semi-leptonic decays of D and B mesons. Some of those results are already available as journal publications: D meson production in pp reactions at $\sqrt{s} = 7$ TeV [11] and in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV [12] and heavy-flavour decays muons in pp interactions at $\sqrt{s} = 7$ TeV [13].

3.1 D mesons at mid-rapidity

D^0 , D^{*+} and D^+ mesons are reconstructed through their hadronic channels: $D^0 \rightarrow K^- \pi^+$ (BR = $(3.91 \pm 0.05)\%$), $D^{*+} \rightarrow D^0 \pi^+$ (BR = $(67.7 \pm 0.5)\%$) and $D^+ \rightarrow K^+ \pi^- \pi^+$ (BR = $(9.22 \pm 0.21)\%$), based on their decay topology and the invariant mass technique. The kaon and pion identification using TPC and TOF helps to reduce background at low p_t . Details of the analysis can be found in [11, 12]. The feed-down from beauty decays is calculated from theory and gives a contribution of 10–15%. The differential cross section for prompt D^0 , D^{*+} and D^+ mesons, measured in the p_t range 1–24 GeV/c, is shown in Fig. 2. The data are well described within uncertainties by perturbative QCD calculations at Fixed-Order plus Next-to-Leading Logarithm (FONLL) level [14] and GM-VFNS predictions [15].

The D meson production cross section was also measured in pp collisions at $\sqrt{s} = 2.76$ TeV ($L_{int} = 1.1 \text{ nb}^{-1}$). The total charm cross section was extracted by extrapolating the measurements to the full phase space. The measured charm cross sections fit into the world data trend and is well described by NLO MNR calculation [16] as shown in Fig. 3. These data are important baseline measurements for quarkonia studies in Pb–Pb collisions and the corresponding parton spectra from pQCD serve as an input for energy loss models.

The first measurement of the nuclear modification factor of prompt D^0 , D^{*+} and D^+ mesons in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV is shown in Fig. 4 for the centrality classes 0–20% and 40–80%. The R_{AA} of the three D meson species agree within statistical uncertainties and shows for the 20% most central collisions a strong suppression by a factor of 4–5 for $p_t > 5$ GeV/c. Less suppression is observed for the centrality class 40–80%. The average R_{AA} of D mesons in the 0–20% central-

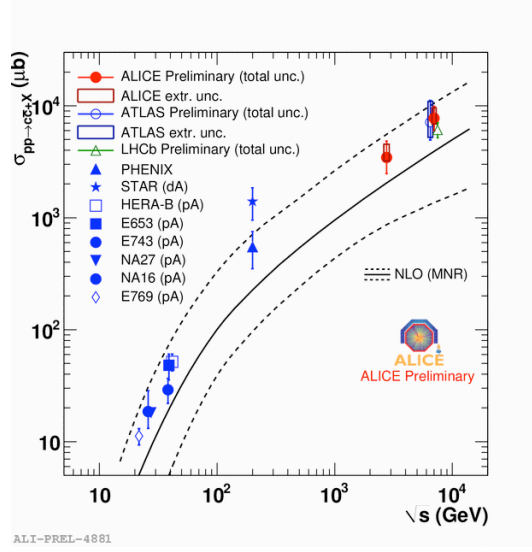


Figure 3: Comparison of total charm cross section in pp collisions at $\sqrt{s} = 2.76$ and 7 TeV with the world data. The NLO MNR calculation [16] (and its uncertainties) is represented by solid (dashed) lines.

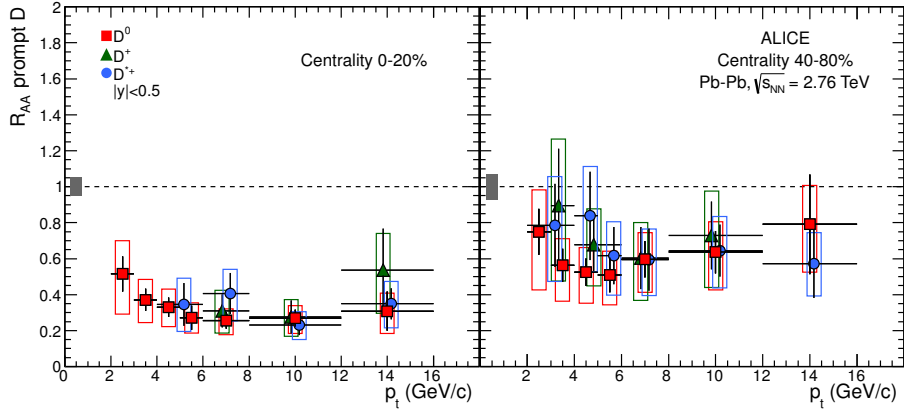


Figure 4: (colour online) R_{AA} for prompt D^0 , D^{*+} and D^+ in the 0–20% (left) and 40–80% (right) centrality classes [12]. Statistical (bars), systematic (empty boxes), and normalization (full box) uncertainties are shown. Horizontal error bars reflect bin widths, symbols were placed at the centre of the bin.

ity class is compatible, within uncertainties, with the R_{AA} of charged hadrons [17] (cf. Fig. 5, left panel). There are indications that $R_{AA}^D > R_{AA}^{h^+}$ at lower p_t , although more data are necessary to draw final conclusions. In Fig. 5 (right panel) the R_{AA}^D measurement is compared with expectation from the NLO pQCD calculations [16] with nuclear shadowing using the EPS09 parametrisation [18]. Only little shadowing is expected in this p_t range. Thus the observed suppression is a final state effect, namely due to the interaction of the heavy quarks with the hot QCD medium. In Fig. 6 the nuclear modification factor for the average D mesons and charged hadrons in the 20% most central Pb-Pb collisions are confronted with several theoretical models based on parton energy loss (see [12] for details). Among the models that compute both observables, radiative energy loss supplemented with in-medium D meson dissociation and radiative plus collisional energy loss in the WHDG and CUJET1.0 implementations simultaneously describe the charm and light-flavour sup-

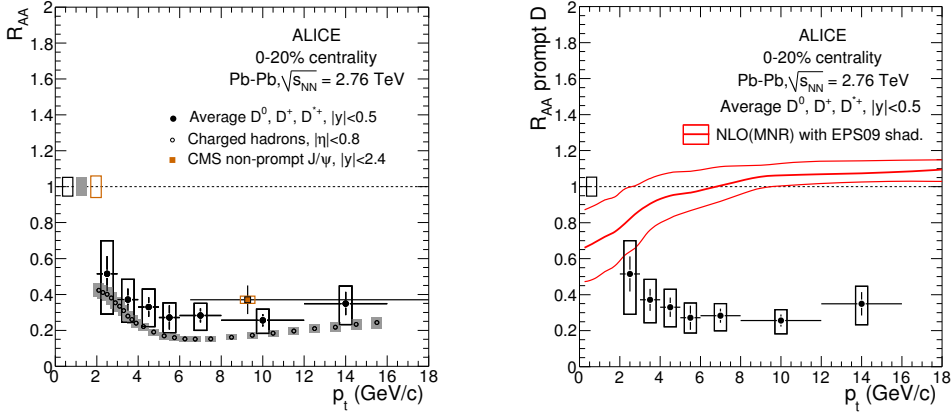


Figure 5: Average R_{AA} of D mesons in the 0–20% centrality class [12] compared to: left, the nuclear modification factors of charged hadrons [17] in the same centrality class and non-prompt J/ψ from B decays; right, the expectation from NLO pQCD [16] with nuclear shadowing [18].

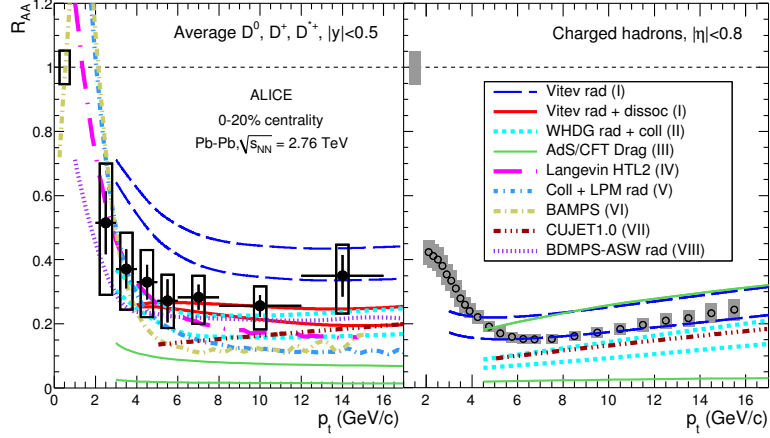


Figure 6: (colour online) Average R_{AA} of D mesons [12] (left) and R_{AA} of charged hadrons [17] (right) in the 0–20% centrality class compared to different model calculations based on parton energy loss (see [12] for details). The two normalisation uncertainties are almost fully correlated.

pression reasonably well. While in the former calculation the medium density is tuned to describe the inclusive jet suppression at LHC energies, for the latter two it is extrapolated to LHC conditions starting from the value that describes the pion suppression at RHIC energy. This might explain why these two models are somewhat low with respect to the charged-hadron R_{AA} . A model based on AdS/CFT drag coefficients underestimates significantly the charm R_{AA} and has very limited predictive power for the light-flavour R_{AA} .

3.2 Single electrons from heavy-flavour decays at mid-rapidity

The cross section for open heavy-flavour production at mid-rapidity has also been studied through the measurement of single electrons. Those electrons are identified, using the TPC and TOF detectors, on a statistical basis by subtracting a cocktail of background electrons from the inclusive

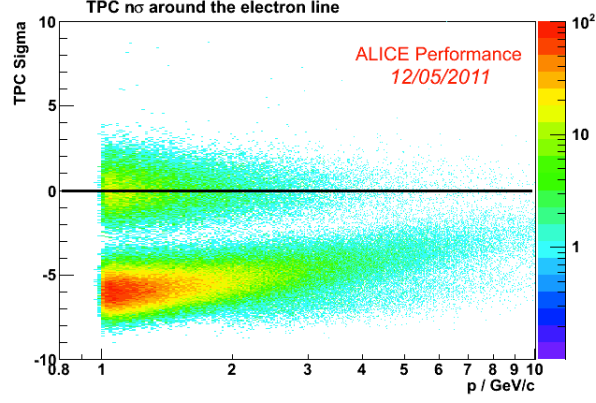


Figure 7: Electron identification using the normalized TPC dE/dx distribution. The upper band is electrons, the lower one charged pions.

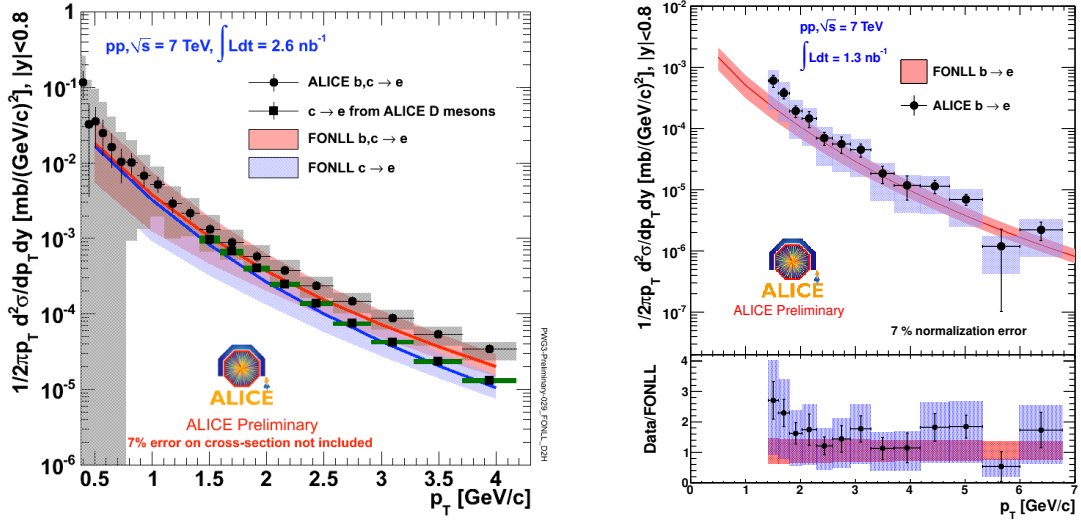


Figure 8: Invariant cross section of electrons from charm or beauty decays (left panel) and beauty decays only (right panel) in pp collisions at $\sqrt{s} = 7$ TeV, compared to FONLL calculations (red and blues curves) and to expectations from D meson decay electrons (left panel).

electron spectrum. This background arises mainly from γ conversion electrons in the detector material and π^0 Dalitz decays. For p_T up to a few GeV/c this cocktail can be determined precisely by means of the measured π^0 cross section. Figure 7 depicts the momentum dependence of the normalized TPC dE/dx distribution after applying a cut on the TOF signal, which rejects kaons (< 1.5 GeV/c) and protons (< 3 GeV/c). A dE/dx cut separates electrons from charged pions up to about 10 GeV/c with a residual pion contamination of less than 15%. Figure 8 (left panel) illustrates the single electron cross section in pp reactions at $\sqrt{s} = 7$ TeV, which has a total systematic uncertainty of 16-20% (p_T dependent) plus 7% for the normalisation. The data are well described by FONLL calculations [14] within uncertainties. Moreover, the low p_T single electron spectrum agrees with expectations from D meson decay electrons. Similar measurements at high- p_T using the EMCAL and TRD detectors are ongoing.

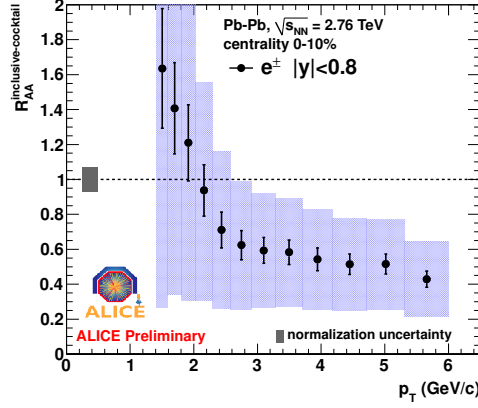


Figure 9: Nuclear modification factor for cocktail-subtracted electrons in the 10% most central Pb-Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV. The blue boxes are the systematic uncertainties.

Electrons from beauty decays are identified through displaced vertices. A selection is applied by requiring electrons to miss the interaction vertex by at least three times the error on the separation of the track from the vertex (of the order of $80\mu\text{m}$ at $2\text{ GeV}/c$). The remaining contribution from charm decays is estimated from the D meson cross section at mid-rapidity and subtracted. The resulting spectrum of electrons from beauty decays is shown in Fig. 8 (right panel) and is well described by FONLL calculations [14] within uncertainties.

A similar analysis was performed with the 2010 Pb–Pb data set. The hadron contamination in the electron sample remains below 10% in the momentum range between 1.5 and $6\text{ GeV}/c$. The nuclear modification factor of cocktail-subtracted electrons for the most 10% central Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV is shown in Fig. 9. The R_{AA}^{electron} and R_{AA}^D at high p_t (cf. Fig. 4) are consistent within uncertainties. Studies are ongoing to determine the bottom contribution to the single electrons experimentally using the electron displacement method and azimuthal angular correlations [19]. However, FONLL calculations indicate that b-hadron decays start to dominate above $\approx 5\text{--}6\text{ GeV}/c$. Thus, beauty suppression appears to be large.

3.3 Single muons from heavy-flavour decays at forward rapidity

Heavy-flavour production at forward rapidities ($-4 < \eta < -2.5$) is studied with single muons using the ALICE muon spectrometer. The extraction of the heavy-flavour contribution of the single muon spectra requires the subtraction of three main background sources: a) muons from the decay-in-flight of light hadrons (decay muons); b) muons from the decay of hadrons produced in the interaction with the front absorber (secondary muons); c) punch-through hadrons. The last contribution can be efficiently rejected by requiring the matching of the reconstructed tracks with the tracks in the trigger system. Due to the lower mass of the parent particles, the background muons have a softer p_t distribution than the heavy-flavour muons, and dominate the low- p_t region. Therefore, the analysis is restricted to the p_t range $2\text{--}12\text{ GeV}/c$. Simulation studies indicate that, in this p_t range, the contribution of secondary muons is small (about 3%). The main source of background in this region consists of decay muons (about 25%), which have been subtracted using Monte Carlo simulations. Figure 10 illustrates the cross section of heavy-flavour decay muons as a function of

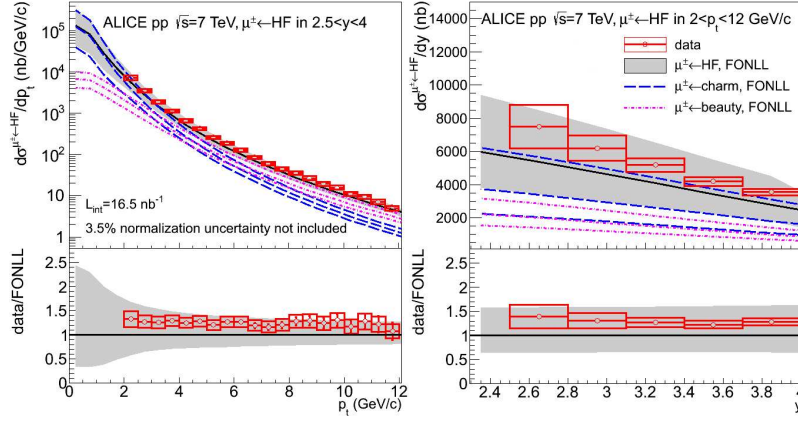


Figure 10: Differential transverse momentum cross section for single muons in $-4 < \eta < -2.5$ [13]. The statistical error is smaller than the markers. The systematic errors (open boxes) do not include an additional 10% error on the minimum-bias pp cross section. The grey band indicates the FONLL prediction.

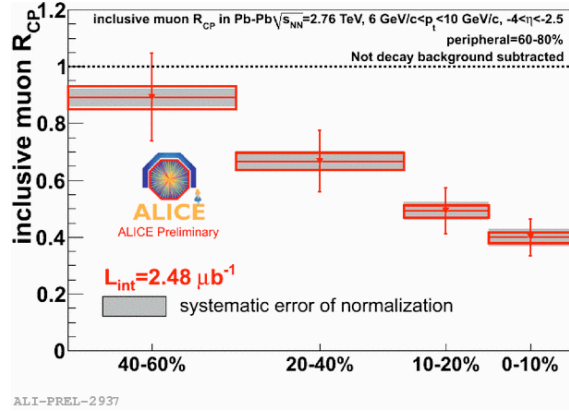


Figure 11: Centrality dependence of the R_{CP} for single muons in the p_t range 6-10 GeV/c at forward rapidity in Pb-Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV.

p_t (left panel) and y (right panel) in pp collisions at $\sqrt{s} = 7$ TeV [13]. The FONLL calculation [14] agrees with the data within uncertainties. Single muons were measured in Pb-Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV using a similar analysis procedure. The R_{CP} for single muons in the p_t range 6-10 GeV/c shows a clear centrality dependence (cf. Figure 11).

4. Summary

Heavy quarks are particularly good probes to study the properties of hot QCD matter (especially the transport properties). The R_{AA} of D^0 , D^{*+} and D^+ was measured for the first time as a function of transverse momentum and centrality in Pb-Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV and indicate strong in-medium energy loss for charm quarks. The suppression is almost as large as that observed for charged (light-flavour) hadrons, with a possible indication of $R_{AA}^D > R_{AA}^{ph\pm}$, although not fully significant with the present level of experimental uncertainties. The expected effect of nuclear shadowing is small ($< 15\%$) above $p_t = 6$ GeV/c, indicating that the large measured suppression cannot be explained by initial-state effects. Some of the pQCD models based on various implementations of

parton energy loss succeed reasonably well at describing simultaneously the suppression of light flavour and charm hadrons. Single electrons and muons show similar suppression at and above $p_t > 5$ GeV/c, where beauty decays are dominant according to NLO pQCD calculations. The precision of the D meson measurements will be improved in the future and the direct measurement of the R_{AA} of electrons from beauty decays should be possible using the large sample of Pb–Pb collisions recorded in 2011. In addition, the upcoming p–Pb run in 2012 will provide insight on cold nuclear matter effects in the low-momentum region.

Acknowledgments

I would like to thank the organisers for inviting me and the stimulating atmosphere at this workshop. Bormio 2012 was a great experience for me.

The European Research Council has provided financial support under the European Community's Seventh Framework Programme (FP7/2007-2013) / ERC grant agreement no 210223. This work was also supported by a Vidi grant from the Netherlands Organisation for Scientific Research (project number 680-47-232) and a Projectruimte grant from the Dutch Foundation for Fundamental Research (project number 10PR2884).

References

- [1] Quark-Gluon Plasma 4, Eds. R.C. Hwa and X.-N. Wang, World Scientific Publishing, 2010.
- [2] Y.L. Dokshitzer, D.E. Kharzeev, Phys. Lett. **B519**, 199 (2001).
- [3] M. Djordjevic, M. Gyulassy, S. Wicks, Phys. Rev. Lett. **94**, 112301 (2005).
- [4] S. Wicks, W. Horowitz, M. Djordjevic, M. Gyulassy, Nucl. Phys. A784, 426 (2007).
- [5] P. Cortese *et al.* [ALICE Collaboration], J. Phys. G: Nucl. Part. Phys. **32**, 1295 (2006).
- [6] K. Aamodt *et al.* [ALICE Collaboration], J. Instrum. **3**, S08002 (2008).
- [7] J. Schukraft, arXiv:1112.0550.
- [8] J.W. Harris (ALICE Collab.), Published by AIP in the Conf. Proc. Ser. (2011), arXiv:1111.4651.
- [9] K. Aamodt *et al.* [ALICE Collaboration], J. Instrum. **5**, P03003 (2010).
- [10] K. Aamodt *et al.*, Phys. Rev. Lett. **105**, 252301 (2010).
- [11] B. Abelev *et al.* (ALICE Collaboration), JHEP **01**, 128 (2012).
- [12] B. Abelev *et al.* (ALICE Collaboration), submitted to JHEP (arXiv:1203.2160), March 2012.
- [13] B. Abelev *et al.* (ALICE Collaboration), Phys. Lett. **B708**, 265 (2012).
- [14] M. Cacciari *et al.*, JHEP **0103**, 006 (2001).
- [15] B.A. Kniehl, G. Kramer, I. Schienbein and H. Spiesberger, arXiv:1202.0439. B.A. Kniehl *et al.*, Phys. Rev. Lett. **96**, 012001 (2006).
- [16] M. Mangano, P. Nason, G. Ridolfi, Nucl. Phys. B **373**, 295 (1992).
- [17] H. Appleshauser, J. Phys. G: Nucl. Part. Phys. **38**, 124014 (2011).
- [18] K.J. Eskola, H. Paukkunen and C.A. Salgado, JHEP **0904**, 065 (2009).
- [19] D. Thomas, Hard Probes 2010. A. Mischke, J. Phys.: Conf. Ser. **316**, 012009 (2011). A. Mischke, Phys. Lett. **B671**, 361 (2009).